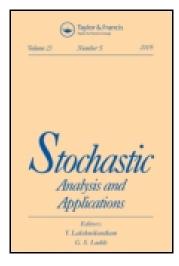
This article was downloaded by: [University of Western Ontario]

On: 24 October 2014, At: 17:26

Publisher: Taylor & Francis

Informa Ltd Registered in England and Wales Registered Number: 1072954 Registered office: Mortimer House, 37-41 Mortimer Street, London W1T 3JH,

UK



Stochastic Analysis and Applications

Publication details, including instructions for authors and subscription information: http://www.tandfonline.com/loi/lsaa20

Numerical Analysis of Stochastic Differential Equations with Explosions

Juan Dávila ^a , Julian Fern ´ndez Bonder ^b , Julio D. Rossi ^b , Pablo Groisman ^c & Mariela Sued ^c ^a Departamento de Ingeniería Matemática , Universidad de Chile , Santiago, Chile ^b Departamento de Matemática , FCEyN, Universidad de Buenos Aires , Buenos Aires, Argentina ^c Instituto de Cálculo , FCEyN, Universidad de Buenos Aires , Buenos Aires , Argentina

To cite this article: Juan Dávila, Julian Fern´ndez Bonder, Julio D. Rossi, Pablo Groisman & Mariela Sued (2005) Numerical Analysis of Stochastic Differential Equations with Explosions, Stochastic Analysis and Applications, 23:4, 809-825, DOI: 10.1081/SAP-200064484

Published online: 15 Feb 2007.

To link to this article: http://dx.doi.org/10.1081/SAP-200064484

PLEASE SCROLL DOWN FOR ARTICLE

Taylor & Francis makes every effort to ensure the accuracy of all the information (the "Content") contained in the publications on our platform. However, Taylor & Francis, our agents, and our licensors make no representations or warranties whatsoever as to the accuracy, completeness, or suitability for any purpose of the Content. Any opinions and views expressed in this publication are the opinions and views of the authors, and are not the views of or endorsed by Taylor & Francis. The accuracy of the Content should not be relied upon and should be independently verified with

primary sources of information. Taylor and Francis shall not be liable for any losses, actions, claims, proceedings, demands, costs, expenses, damages, and other liabilities whatsoever or howsoever caused arising directly or indirectly in connection with, in relation to or arising out of the use of the Content.

This article may be used for research, teaching, and private study purposes. Any substantial or systematic reproduction, redistribution, reselling, loan, sub-licensing, systematic supply, or distribution in any form to anyone is expressly forbidden. Terms & Conditions of access and use can be found at http://www.tandfonline.com/page/terms-and-conditions

Stochastic Analysis and Applications, 23: 809-825, 2005

Copyright © Taylor & Francis, Inc. ISSN 0736-2994 print/1532-9356 online

DOI: 10.1081/SAP-200064484

Numerical Analysis of Stochastic Differential Equations with Explosions

Juan Dávila

Departamento de Ingeniería Matemática, Universidad de Chile, Santiago, Chile

Julian Fernández Bonder and Julio D. Rossi

Departamento de Matemática, FCEyN, Universidad de Buenos Aires, Buenos Aires, Argentina

Pablo Groisman and Mariela Sued

Instituto de Cálculo, FCEyN, Universidad de Buenos Aires, Buenos Aires, Argentina

Abstract: Stochastic ordinary differential equations may have solutions that explode in finite or infinite time. In this article we design an adaptive numerical scheme that reproduces the explosive behavior. The time step is adapted according to the size of the computed solution in such a way that, under adequate hypotheses, the explosion of the solutions is reproduced.

Keywords: Explosion; Numerical approximations; Stochastic differential equations.

Mathematics Subject Classification: 65C30; 65L20; 60H10.

Received December 23, 2004; Accepted February 9, 2005 Address correspondence to Pablo Groisman, Instituto de Cálculo, FCEyN, Universidad de Buenos Aires, Pabellón II, Ciudad Universitaria (1428), Buenos Aires, Argentina; E-mail: pgroisma@dm.uba.ar

1. INTRODUCTION

Consider the following stochastic differential equation (SDE):

$$dx = b(x)dt + \sigma(x)dW, \tag{P}$$

with $x(0) = z \in \mathbb{R}_{>0}$, where b and σ are smooth positive functions (C^1 or even locally Lipschitz will be enough for our calculations) and W is a (one-dimensional) Wiener process defined on a given complete probability space $(\Omega, \mathcal{F}, \mathbb{P})$ with a filtration $\{\mathcal{F}_t\}_{t\geq 0}$ satisfying the usual conditions (i.e., it is right continuous and \mathcal{F}_0 contains all \mathbb{P} -null sets [4]).

It is well known that stochastic differential equations like (P) may explode in finite time. That is, trajectories may diverge to infinity as t goes to some finite time T that, in general, depends on the particular path.

The Feller Test for explosions (see [4, 6]) gives a precise description in terms of b, σ , and z of whether explosions in finite time occur with probability zero, positive, or one. We review some well known facts about SDE with explosions in Section 2.

For example, if b and σ behave like powers at infinity, i.e.,

$$b(s) \sim s^p$$
, $\sigma(s) \sim s^q \ (s \to \infty)$,

applying the Feller test, one obtains that solutions to (P) explode with probability one if $p > 2q \lor 1$. We use $f(s) \sim g(s)$ to mean that there exist constants 0 < c < C, such that $cg(s) \le f(s) \le Cg(s)$ for large enough s. We also use $a \lor b = \max\{a, b\}$, $a \land b = \min\{a, b\}$.

The intuition behind this condition is that p > 2q ensures that the asymptotic behavior of the solutions is governed by the drift term while p > 1 impose the solution to grow up so fast that it explodes in finite time, as in the deterministic case ($\sigma = 0$).

Stochastic differential equations like (P) have been considered, for example, in fatigue cracking (fatigue failures in solid materials) with b and σ of power type [5, 9], and so solutions may explode in finite time. This explosion time is generally random, depends on the particular sample path, and corresponds to the time of ultimate damage or fatigue failure in the material.

Unfortunately, explicit solvable SDEs are rare in practical applications, hence the importance of developing numerical methods to approximate them.

There are many numerical methods designed to deal with SDEs like (P) when b and σ are assumed to be globally Lipschitz continuous. See, for instance, the surveys [2, 7] and the book [5]. See also [3], where locally Lipschitz coefficients are considered. However, all of the cited work deal with globally defined solutions, and most of them with a constant

time step. When dealing with explosive solutions, these methods do not apply mainly because using a constant time step produces approximations that are globally defined. Moreover, the convergence results are based on regularity assumptions of the solution in a fixed (deterministic) time interval $[0, \tau]$; these hypotheses are not available in our case.

The main purpose of this article is to develop an adaptive method that reproduces explosions of the solutions in case that it occurs, providing rigorous proofs of this fact.

We want to remark that even for deterministic problems, the usual numerical methods are not well suited to reproduce explosions, and therefore, adaptive schemes have been developed [1].

The Numerical Scheme

Let h > 0 be the parameter of the method and let $\{X_k\}_{k \ge 1} = \{X_k^h\}_{k \ge 1}$ be the numerical approximation of (P) given by the Euler-Maruyama method

$$X_{k+1} = X_k + \tau_k b(X_k) + \sigma(X_k) \Delta W_k, \quad X_0 = x(0) = z,$$
 (P_h)

where $\Delta W_k = W_{t_{k+1}} - W_{t_k}$ denotes the increment of the Wiener process in the interval $[t_k, t_{k+1}]$ and $\tau_k = t_{k+1} - t_k$.

We define, for notational purposes, X(t) as an interpolant of $X(t_k) = X_k$. For instance, we can take $X(t) \equiv X_k$ for $t_k \le t < t_{k+1}$ or X(t) to be the linear interpolant of the values X_k .

Observe that the numerical approximation X(t) is a well-defined process up to time

$$T_h := \sum_{k=1}^{\infty} \tau_k.$$

We say that a sample path $X(\cdot, \omega)$ of (P_h) explodes in finite time if

$$X(t, \omega) \to +\infty \quad (t \to T_h), \quad \text{and} \quad T_h(\omega) < \infty.$$

If b and σ are globally Lipschitz continuous, it is customary to take a constant time step $\tau_k \equiv h$ (see [5]). However, when designing adaptive algorithms, the time step τ_k has to be selected according to the computed solution X_k and so it will be necessarily aleatory. Inspired by [1], we select the time steps τ_k , according to the rule

$$\tau_k = \frac{h}{b(X_k)}.$$

Observe that by our selection of τ_k , $T_h(\omega) < \infty$ implies $X(t) \nearrow +\infty$.

Main Results

First, we prove convergence of the numerical approximations in compact (random) intervals where the solution and the numerical approximation are bounded. For this theorem the time steps τ_k only need to be \mathcal{F}_{t_k} -measurable and verify $\tau_k \leq Ch$, but are otherwise arbitrary.

Theorem 1.1 (Convergence of the numerical scheme). Let $x(\cdot)$ be the solution of (P) and $X(\cdot)$ its EM approximation given by (P_h) . Fix a time $\bar{\tau} > 0$ and a constant M > 0. Consider the stopping times given by $\tau := \bar{\tau} \land R_h^M$ and $\tau_h := \bar{\tau} \land R_h^{2M}$, where $R^M := \inf\{t : |x(t)| \ge M\}$ and $R_h^M := \inf\{t : |X(t)| \ge M\}$. Then

$$\lim_{h\to 0} \mathbb{E}\Big[\sup_{0\le t\le \tau\wedge\tau_h} |x(t)-X(t)|^2\Big] = 0.$$

Observe that, if the sample paths are uniformly bounded, this is a standard convergence theorem (see, for example, [3, 5]). However, in case that there exists solutions that explode in finite time, we prove convergence of the numerical scheme in regions where they are bounded. We do not expect convergence in bigger regions.

However, if we weaken the notion of convergence, we can prove that the computed solution converges to the continuous one in any interval where the continuous solution remains bounded. More precisely, we have

Corollary 1.2. With the same assumptions and notation as in Theorem 1.1, for every $\varepsilon > 0$ and every $0 \le \alpha < \frac{1}{2}$,

$$\mathbb{P}\Big(\sup_{0 \le t \le R^M} |x(t) - X(t)| > \varepsilon h^{\alpha}\Big) \to 0 \quad as \ h \to 0.$$

Next we assume a specific behavior on the coefficients in (P) to have explosions with probability one. The precise assumptions on b, σ are: There exist positive constants κ_1 , κ_2 such that

$$\kappa_1 \le \sigma^2(s) \le \kappa_2 b(s), \quad b \text{ is nondecreasing and } \int_0^\infty \frac{1}{b(s)} \, ds < \infty.$$
(A)

Remark 1.1. By means of the Feller Test, one can check that under assumption (A) solutions to (P) explode with probability one. These assumptions are actually stronger than the ones required by the Feller Test. Recall that the solution of the (deterministic) differential equation $\bar{x}'(t) = b(\bar{x}(t))$ with $\bar{x}(0) = z \in \mathbb{R}_{>0}$ explodes in finite time, if and only if, $\int_{z}^{+\infty} \frac{1}{b(s)} ds < \infty$.

Next, we analyze the asymptotic behavior of the solutions to (P_h) and show that it agrees with the behavior of the solutions to (P). This is our main result.

Theorem 1.3. Assume (A). Then

- 1. For every initial datum z > 0, $X(\cdot)$ explodes in finite time with probability one.
- 2. We have,

$$\lim_{k \to \infty} \frac{X(t_k)}{hk} = 1 \quad a.s.$$

Moreover, for any $\alpha > 1$, there exits $k_0 = k_0(\omega)$, such that, for every $k \ge k_0$, there holds

$$\sum_{j=k}^{\infty} \frac{h}{b(\alpha h j)} \le T_h - t_k \le \sum_{j=k}^{\infty} \frac{h}{b(\alpha^{-1} h j)}.$$

3. If b has regular variation at infinity (see Definition 1),

$$\lim_{k \to \infty} \frac{T_h - t_k}{\int_{X_h}^{\infty} \frac{1}{b(s)} ds} = 1 \quad a.s.$$

4. In addition, for every h > 0, there holds $h/b(z) \le T_h < +\infty$ a.s and for every L > 0, $\mathbb{P}(T_h > L) > 0$.

Remark 1.2. Observe that 3, gives the precise asymptotic behavior of the numerical solution near the explosion time. For example, if $b(s) \sim s^p$, the explosion rate given by 3 is

$$X(t_k)(T_h - t_k)^{1/(p-1)} \to \left(\frac{1}{p-1}\right)^{\frac{1}{p-1}}, \quad (t_k \to T_h).$$

This is the behavior of solutions to the deterministic ODE $d\bar{x}(t) = \bar{x}^p(t)dt$.

Finally, we analyze the convergence of the stopping times considered in Theorem 1.1 to the explosion time of the continuous problem.

Theorem 1.4. Assume (A). Then, for any $\varepsilon > 0$,

$$\lim_{M\to+\infty}\lim_{h\to 0}\mathbb{P}(|R_h^M-T|>\varepsilon)=0.$$

This last theorem is useful in actual computations of the explosion time for (P). See Section 5.

2. THE CONTINUOUS EQUATION

In this section we review some results concerning the behavior of solutions to (P) as $t \nearrow T$, the explosion time. These results can be found, for instance, in [4, 8].

Let $s: \mathbb{R} \to \mathbb{R}$ be the scale function for (P) given by

$$s(z) = 0,$$
 $s'(\xi) = \exp\left[-\int_0^{\xi} 2b(t)\sigma(t)^{-2}dt\right].$

Then, if y(t) = s(x(t)), we have

$$dy = \bar{\sigma}(y)dW,\tag{1}$$

where $\bar{\sigma} = (s'\sigma) \circ s^{-1}$. Solutions to (1) are globally defined. Observe that x explodes in finite time, if and only if

$$\ell := s(+\infty) < +\infty.$$

We can obtain a weak solution to (1) by time change. In fact, let B(t) be a standard Brownian motion and define

$$A(t) = \int_0^t \bar{\sigma}(B(u))^{-2} du,$$

and let γ be the inverse of A, then

$$y(t) = B(\gamma(t))$$

is a weak solution of (1).

Let

$$T_{\ell} := \inf\{t > 0 : B(t) = \ell\}.$$

Therefore

$$T = A(T_{\ell}) = \int_{0}^{T_{\ell}} \bar{\sigma}^{-2}(B(u))du, \tag{2}$$

is the explosion time.

To describe the behavior of x(t) near the explosion time T, we have to study the behavior of B(t) when t is close to T_{ℓ} . To this end, we define

$$R(t) := \ell - B(T_{\ell} - t), \quad 0 \le t \le T_{\ell}.$$

Then R(t) is a Bessel process, i.e., $R(t) \stackrel{\mathcal{L}}{=} BES(3)$.

Combining these assertions, we get

$$y(T - \varepsilon) = B(\gamma(T - \varepsilon)) = B(T_{\ell} - (T_{\ell} - \gamma(T - \varepsilon)))$$
$$= \ell - R(T_{\ell} - \gamma(T - \varepsilon)).$$

So we arrive at

$$x(T-\varepsilon) = s^{-1}(y(T-\varepsilon)) = s^{-1}(\ell - R(T_{\ell} - \gamma(T-\varepsilon))).$$

Therefore, we have found the asymptotic behavior for the solution x to (P) near the explosion time T. Moreover, (2) gives an "explicit" formula for the explosion time T of weak solutions to (P).

3. CONVERGENCE OF THE NUMERICAL SCHEME

We begin this section by showing some measurability properties of the numerical scheme.

Lemma 3.1. With the notation of Section 1.2, $\{t_k\}_{k\geq 1}$ are stopping times and each τ_k is \mathcal{F}_{t_k} -measurable.

Proof. We just observe that

$$t_{k+1} = t_k + \tau_k = t_k + \frac{h}{b(X_k)}.$$

Assume t_k is a stopping time. Then X_k (and hence τ_k) is \mathcal{F}_{t_k} -measurable. Since τ_k is positive, t_{k+1} is also a stopping time. The case k=0 holds since t_0 is the constant h/b(z).

Now we prove the main result of the section. Recall that this result and the subsequent proposition and corollary hold true for any choice of time steps τ_k , if they are \mathcal{F}_{t_k} -measurables and $\tau_k \leq Ch$.

Proof of Theorem 1.1. First, we truncate the functions b(x) and $\sigma(x)$ in such a way that they are globally Lipschitz, bounded, and coincide with the original b(x) and $\sigma(x)$ for values of x with $|x| \le 2M$, i.e., we consider

$$\bar{b}(x) = \begin{cases} b(x) & \text{if } |x| \le 2M \\ b(2M) & \text{if } x \ge 2M \\ b(-2M) & \text{if } x \le -2M, \end{cases}$$

and

$$\bar{\sigma}(x) = \begin{cases} \sigma(x) & \text{if } |x| \le 2M\\ \sigma(2M) & \text{if } x \ge 2M\\ \sigma(-2M) & \text{if } x < -2M. \end{cases}$$

Let y and Y be the solutions of

$$dy = \bar{b}(y)dt + \bar{\sigma}(y)dW, \quad y(0) = z,$$
(3)

$$Y_{k+1} = Y_k + \tau_k \bar{b}(Y_k) + \bar{\sigma}(Y_k) \Delta W_k, \quad Y(0) = z,$$
 (4)

respectively.

From [5], we have

$$\mathbb{E}\Big[\sup_{0 \le t \le \bar{\tau}} |y - Y|^2\Big] \to 0, \text{ as } h \to 0,$$

Recalling that $\tau := \bar{\tau} \wedge R^M$ and $\tau_h := \bar{\tau} \wedge R_h^{2M}$ we have that x(t) = y(t) and X(t) = Y(t) if $0 \le t \le \tau \wedge \tau_h$. Hence

$$\mathbb{E}\Big[\sup_{0 \le t \le \tau \wedge \tau_h} |x - X|^2\Big] = \mathbb{E}\Big[\sup_{0 \le t \le \tau \wedge \tau_h} |y - Y|^2\Big] \le \mathbb{E}\Big[\sup_{0 \le t \le \tilde{\tau}} |y - Y|^2\Big].$$

This implies the result.

Remark 3.1. Observe that, in fact, the results in [5] give

$$\mathbb{E}\Big[\sup_{0\leq t\leq \bar{\tau}}|y-Y|^2\Big]\leq Ch,$$

so, in our case, we also obtain

$$\mathbb{E}\Big[\sup_{0 \le t \le \tau \land \tau_h} |x - X|^2\Big] \le Ch.$$

What one really wants in Theorem 1.1 is convergence of the numerical scheme without any assumptions on the computed solution X. Unfortunately, we are not able to prove convergence in square mean without this hypothesis. However, we are able to prove convergence in probability without any further assumption on X. To this end, we need the following proposition.

Proposition 3.2. Let R^M and R_h^M be as in Theorem 1.1 and M > 0, then we have

- 1. $\mathbb{P}(R^M \ge R_h^{2M}) \to 0$ as $h \to 0$. 2. $\mathbb{P}(R_h^M \ge R^{2M}) \to 0$ as $h \to 0$.

Proof. First, we prove 1. Let $\varepsilon > 0$. As $T < \infty$ a.s., we can take $\bar{\tau}$ such that

$$\mathbb{P}(T>\bar{\tau})<\frac{\varepsilon}{2}.$$

Now, with the notation of the Theorem 1.1, we have that

$$\mathbb{P}(R^{M} \ge R_{h}^{2M}) \le \mathbb{P}(T > \bar{\tau}) + \mathbb{P}\left(\sup_{0 < t < \tau} Y \ge 2M\right)
< \frac{\varepsilon}{2} + \mathbb{P}\left(\sup_{0 < t < \tau} Y \ge 2M\right)
< \varepsilon,$$
(5)

if h is small enough. In fact, by Tchebychev inequality,

$$\mathbb{P}\Big(\sup_{0\leq t\leq \tau}Y\geq 2M\Big)\leq \mathbb{P}\Big(\sup_{0\leq t\leq \tau}|y-Y|>M\Big)\leq \frac{1}{M^2}\mathbb{E}\Big[\sup_{0\leq t\leq \tilde{\tau}}|y-Y|^2\Big]<\frac{\varepsilon}{2},$$

from where 1 follows.

To prove 2 taking $\bar{\tau}$ as before, we have

$$\mathbb{P}(R_h^M \ge R^{2M}) \le \mathbb{P}(T \ge \bar{\tau}) + \mathbb{P}(R_h^M \ge R^{2M}, \ T < \bar{\tau})$$
$$\le \frac{\varepsilon}{2} + \mathbb{P}\Big(\sup_{0 \le t \le \bar{\tau} \land R_h^M} y \ge 2M\Big).$$

The proof follows as in 1.

Now, combining Theorem 1.1 and Proposition 3.2, we can get rid of the boundedness assumption on X by weakening the notion of convergence.

Proof of Corollary 1.2. First, take $\bar{\tau} > 0$ such that

$$\mathbb{P}(T>\bar{\tau})<\frac{\delta}{2}.$$

Then, if $\tau = \bar{\tau} \wedge R^M$, by Tchebychev's inequality,

$$\begin{split} & \mathbb{P}\Big(\sup_{0 < t < \tau} |x(t) - X(t)| > \varepsilon h^{\alpha}\Big) \\ & \leq \mathbb{P}\Big(\sup_{0 < t < \tau} |x(t) - X(t)| > \varepsilon h^{\alpha}, \ R^{M} < R_{h}^{2M}\Big) + \mathbb{P}(R^{M} \ge R_{h}^{2M}) \\ & \leq \frac{1}{h^{2\alpha}\varepsilon^{2}} \mathbb{E}\Big[\sup_{0 < t < \tau \wedge \tau_{h}} |x(t) - X(t)|^{2}\Big] + \mathbb{P}(R^{M} \ge R_{h}^{2M}) \\ & \leq \frac{h^{1-2\alpha}}{\varepsilon} + \mathbb{P}(R^{M} \ge R_{h}^{2M}) \to 0 \end{split}$$

as $h \to 0$. Therefore, for h small (depending on δ)

$$\mathbb{P}\Big(\sup_{0 < t < R^M} |x(t) - X(t)| > \varepsilon h^{\alpha}\Big) < \delta.$$

This finishes the proof.

4. EXPLOSIONS IN THE NUMERICAL SCHEME

In this section, we prove that for almost every $\omega \in \Omega$, X(t) explodes in finite time $T_h(\omega)$.

To begin with, let us recall an auxiliary lemma.

Lemma 4.1 ([4], Chapter 2). Let $\{\mathcal{F}_t\}$ be the filtration generated by the Wiener process $W(\cdot)$ and S a stopping time of $\{\mathcal{F}_t\}$. Assume τ is a random variable \mathcal{F}_S -measurable. Then, for every Borel set A, it holds

$$\mathbb{P}(W(S+\tau)-W(S)\in A|\mathcal{F}_S) = \int_A \frac{1}{\sqrt{2\pi\tau}} e^{-\frac{x^2}{2\tau}} dx. \tag{6}$$

Furthermore,

$$\frac{W(S+\tau)-W(S)}{\sqrt{\tau}}\Big|_{\mathscr{F}_S}$$

is a standard normal random variable and hence (6) holds without conditioning.

Next, we prove a technical lemma that is the key point in the proof of Theorem 1.3. This lemma allows us to control the effect of the diffusion in the numerical approximations of (P).

Lemma 4.2. Let $Y_k = \sum_{i=1}^k \sigma(X_i) \Delta W_i$. Then

$$\lim_{k \to \infty} \frac{Y_k}{k} = 0 \quad a.s.$$

Proof. Let

$$Z_j := \frac{\Delta W_j}{\sqrt{\tau_j}} = \frac{W_{t_j + \tau_j} - W_{t_j}}{\sqrt{\tau_j}} \quad \text{and} \quad a_j := \sigma(X_j) \sqrt{\tau_j} = \sqrt{h} \frac{\sigma(X_j)}{\sqrt{b(X_j)}}.$$

Then $Y_k = \sum_{j=1}^k a_j Z_j$. Observe that a_j are uniformly bounded by assumption (A). In order to prove that Y_k/k goes to zero, we use Tchebychev's inequality combined with Borel-Cantelli's Lemma. So we need to show that

$$\sum_{k=1}^{\infty} \frac{\mathbb{E}[Y_k^4]}{k^4} < \infty.$$

Observe that Z_j is independent of \mathcal{F}_{t_j} and is normally distributed, according Lemma 4.1. Then, if $i \neq j$, r or l, conditioning we obtain

$$\mathbb{E}[Z_i Z_j Z_r Z_l] = 0.$$

Moreover,

$$\mathbb{E}[Z_i^2 Z_i^2] = 1 \ (i \neq j) \text{ and } \mathbb{E}[Z_i^4] = 3.$$

Hence,

$$\mathbb{E}\left[\left(\sum_{j=1}^{k} Z_{j}\right)^{4}\right] = \sum_{j=1}^{k} \mathbb{E}[Z_{j}^{4}] + 3 \sum_{\substack{i,j=1\\i\neq j}}^{k} \mathbb{E}[Z_{i}^{2}Z_{j}^{2}]$$
$$= 3k + 3(k^{2} - k)$$
$$= 3k^{2}.$$

Taking into account that a_j is \mathcal{F}_{t_j} -measurable, proceeding in the same way with $a_i Z_j$, we obtain

$$\mathbb{E}[Y_k^4] \le 3(\kappa_2 h k)^2,$$

to get the desired result.

Now we use this lemma to prove that solutions to the numerical scheme explode with probability one and to find the rate of explosion. We are going to use the following

Definition 1. We say that a function $f: \mathbb{R} \to \mathbb{R}$ has regular variation at infinity if there exist p > 0, such that for every positive α ,

$$\lim_{s\to+\infty}\frac{f(\alpha s)}{f(s)}=\alpha^p.$$

Proof of Theorem 1.3. Since $X_k = z + hk + Y_{k-1}$, by Lemma 4.2,

$$\frac{X_k}{hk} = \frac{z}{hk} + 1 + \frac{Y_{k-1}}{hk} \to 1, \quad \text{a.s. as } k \to \infty.$$
 (7)

To prove that explosion occurs, it remains to be shown that with probability one, $T_h = \sum \tau_j < \infty$. To this end observe that, by (7), for every $\alpha > 1$, there exist $k_0 = k_0(\omega)$, such that,

$$\begin{split} \sum_{k=k_0}^{\infty} \tau_k &= \sum_{k=k_0}^{\infty} \frac{h}{b(X_k)} \le \sum_{k=k_0}^{\infty} \frac{h}{b(\alpha^{-1}hk)} \le \int_{k_0-1}^{\infty} \frac{h}{b(\alpha^{-1}hs)} \, ds \\ &= \alpha \int_{\alpha^{-1}h(k_0-1)}^{\infty} \frac{1}{b(u)} \, du \le \alpha \int_{X_{k_0-1}}^{\infty} \frac{1}{b(u)} \, du < +\infty, \quad \text{a.s.} \end{split}$$

This proves 1.

In order to prove 2, we observe that the computations performed give

$$T_h - t_k \le \sum_{j=k}^{\infty} \frac{h}{b(\alpha^{-1}hj)} \le \alpha \int_{X_k}^{\infty} \frac{1}{b(u)} du.$$

In the same way, we can obtain the reverse inequality,

$$T_h - t_k \ge \sum_{j=k}^{\infty} \frac{h}{b(\alpha h j)} \ge \frac{1}{\alpha} \int_{X_{k+1}}^{\infty} \frac{1}{b(u)} du,$$

for $k = k(\omega)$ large enough.

To prove 3, just observe that regular variation at infinity of b implies that

$$B(t) := \int_{t}^{\infty} \frac{1}{b(u)} du$$

has regular variation at infinity with the same exponent. Therefore, as *B* is increasing,

$$\frac{B(X_k)}{B(X_{k+1})} = \frac{B(\frac{X_k}{X_{k+1}}X_{k+1})}{B(X_{k+1})} \le \frac{B(\alpha X_{k+1})}{B(X_{k+1})} \to \alpha^p \quad \text{as } k \to \infty$$

for any $\alpha > 1$. Therefore

$$\limsup_{k\to\infty}\frac{B(X_k)}{B(X_{k+1})}\leq 1.$$

Analogously,

$$\liminf_{k\to\infty}\frac{B(X_k)}{B(X_{k+1})}\geq 1.$$

It remains to show 4, but this follows from the fact that for any K > 0,

$$\mathbb{P}\Big(\max_{1\leq j\leq K}X_j<2z\Big)>0.$$

Hence, if K is such that Kh/b(2z) > L, we obtain

$$\mathbb{P}(T_h > L) \ge \mathbb{P}\Big(\max_{1 \le j \le K} X_j < 2z\Big) > 0.$$

Moreover, $T_h \ge \tau_1 = h/b(z)$. The proof is now complete.

5. APPROXIMATION OF THE EXPLOSION TIME

In this section we prove Theorem 1.4. Observe that in numerical simulations R_h^M can be easily computed. This fact together with

Theorem 1.4 allows us to construct the numerical approximation of T given in the next section.

Proof of Theorem 1.4. We proceed as follows, for $\varepsilon > 0$, we have

$$\mathbb{P}(|R_h^M - T| > \varepsilon) = \mathbb{P}(R_h^M - T > \varepsilon) + \mathbb{P}(R_h^M - T < -\varepsilon) = I + II.$$

We first show that I goes to zero as $h \to 0$ for any fixed M. In fact, as $R^{2M} < T$,

$$I \leq \mathbb{P}(R_h^M - R^{2M} > \varepsilon) \leq \mathbb{P}(R_h^M > R^{2M}) \to 0,$$

as $h \to 0$, by Proposition 3.2.

For the second term, we have

$$II \le \mathbb{P}(R^{M/2} - R_h^M > \varepsilon/2) + \mathbb{P}(T - R^{M/2} > \varepsilon/2)$$

$$\le \mathbb{P}(R^{M/2} > R_h^M) + \mathbb{P}(|T - R^{M/2}| > \varepsilon/2) \to 0$$

by Proposition 3.2 and since $R^{M/2} \to T$ a.s. $(M \to +\infty)$. This completes the proof.

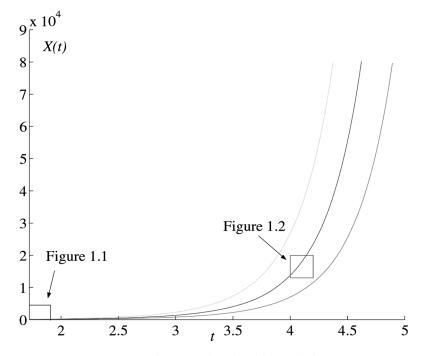


Figure 1. Three sample paths with explosions.

6. NUMERICAL EXPERIMENTS

In this section, we present some numerical examples to illustrate the theory set forth in the previous sections. All the experiments are computed with

$$b(\xi) = |\xi|^{1.1} + 0.1, \qquad \sigma(\xi) = \sqrt{|\xi|^{1.1} + 0.1}, \qquad z = 10.$$

The increments of the Wiener process have been generated with the randn routine of MATLAB.

In Figure 1, we show some sample paths of the solution. We stop the algorithm when the computed solution reaches $M = 10^5$.

In Figure 2, we show the ratio X_k/hk , and observe that it converges to 1 as predicted by our result.

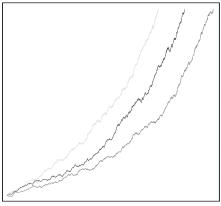


Figure 1.1.

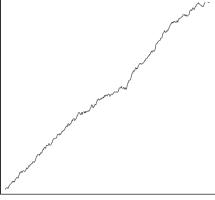


Figure 1.2.

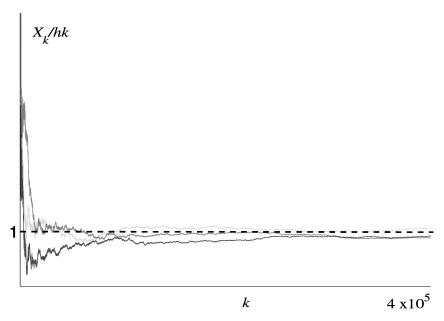


Figure 2. $X_k/hk \rightarrow 1$ a.s.

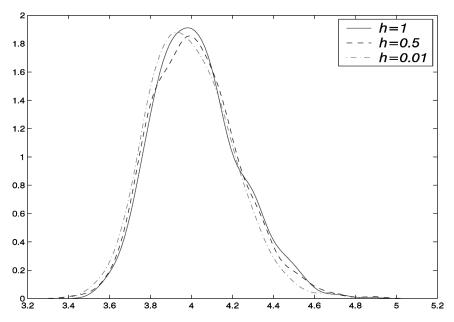


Figure 3. The kernel estimator of the densitry of R_h^M for different values of h.

Finally, Figure 3 shows the kernel density estimation of R_h^M for different choices of h. As proven in Theorem 1.4, R_h^M converges in probability to T.

We have used 1000 sample paths for each estimator. The values of h taken in each estimator were h=1, h=0.5, and h=0.1. Observe that, in each case, the largest time step taken was $\tau_1 \simeq 0.08$, $\tau_1 \simeq 0.04$, and $\tau_1 \simeq 0.008$, respectively.

ACKNOWLEDGMENTS

We like to thank Jean Bertoin and Pablo Ferrari for several interesting discussions that helped us improve this article. J.F.B and P.G. are grateful for the hospitality received at the University of Sao Paulo, where part of the work was done. J.D. is grateful for the invitation and hospitality at the Department of Mathematics at the University of Buenos Aires. This work is supported by Universidad de Buenos Aires under grant TX066, by ANPCyT PICT No. 03-05009, Fundación Antorchas Project 13900-5 and CONICET (Argentina). J.F.B. is also supported by ANPCyT PICT No. 03-10608. The manuscript is supported by a fellowship of Fundacion Antorchas.

REFERENCES

- 1. Acosta, G., R. Durán, and J. D. Rossi. 2002. An adaptive time step procedure for a parabolic problem with blow-up. *Computing* 68(4):343–373.
- Higham, D. J. 2001. An algorithmic introduction to numerical simulation of stochastic differential equations. SIAM Rev. 43(3): 525–546.
- 3. Higham, D. J., X. Mao, and A. M. Stuart. 2002. Strong convergence of Euler-type methods for nonlinear stochastic differential equations. *SIAM J. Number. Anal.* 40(3):1041–1063.
- 4. Karatzas, I., and S. E. Shreve. 1991. Brownian motion and stochastic calculus, volume 113 of Graduate Texts in Mathematics, 2nd ed. New Work: Springer-Verlag.
- Kloeden, P. E., and E. Platen. 1992. Numerical Solution of Stochastic Differential Equations, Volume 23 of Applications of Mathematics (New York). Berlin: Springer-Verlag.
- 6. McKean, H. P. Jr. 1969. Stochastic Integrals. *Probability and Mathematical Statistics*, No. 5. New York: Academic Press.
- 7. Platen, E. 1999. An introduction to numerical methods for stochastic differential equations. In *Acta numerica*, Volume 8 of Acta Numer., 197–246. Cambridge: Cambridge University Press.

- 8. Rogers, L. C. G., and D. Williams. 1994. Diffusions, Markov Processes, and Martingales. Vol. 2. Cambridge Mathematical Library. Cambridge: Cambridge University Press.
- 9. Sobczyk, K., and B. F. Spencer, Jr. 1992. Random Fatigue. Boston, MA: Academic Press Inc.